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Noise-Induced Quantum Synchronization and Entanglement in a Quantum Analogue of Huygens' Clock

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08/2024

Asymmetry Amplification by a Nonadiabatic Passage through a Critical Point

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arXiv: 2408.15897

PRA: In Review



10/2024

Noise-induced synchronization in coupled quantum oscillators

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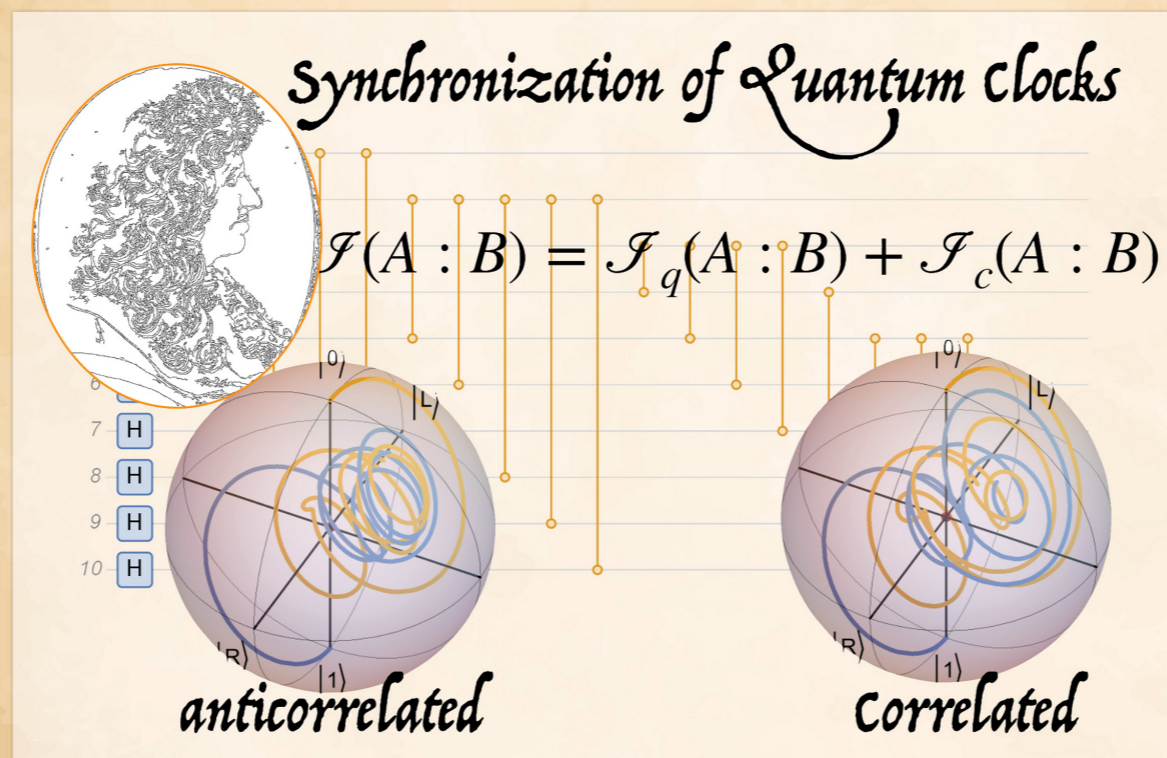
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Bhavay Tyagi
Advisor: Prof. Eric R. Bittner
Theoretical Chemical Physics Group

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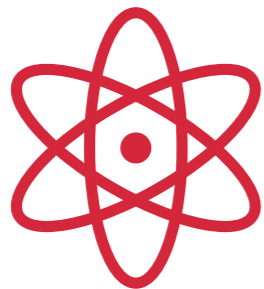
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A **Big** (sometimes solvable) Problem

- The Problem: **Decoherence**
 - Thermal Fluctuations
 - Nuclear, Electronic and Vibrational Fluctuations
Even in the most ideal case
 - Vacuum and Fields Fluctuations
- The Want: To make our system **robust** against this or at least control it.



$$\tau_D = 10^{-6} s$$



$$\tau_D = 10^{-30} s$$



$$\tau_P = 10^{-3} s$$

$$\tau_P = 100 ms$$

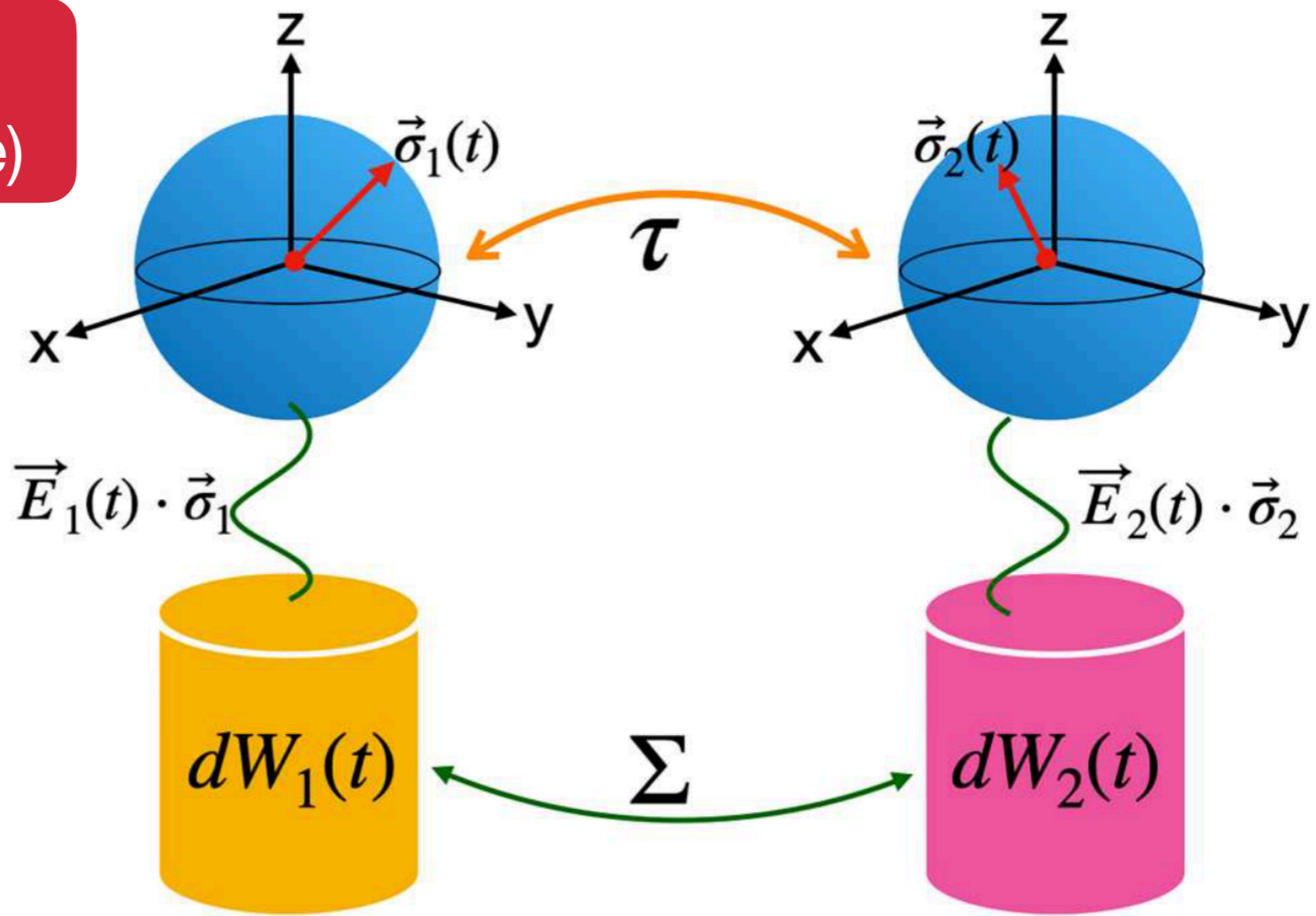
A **Big** (sometimes solvable) Problem

- The Problem: **Decoherence**
- The Want: To make our system **robust** against this or at least control it.
- The **Big** Idea: What if this **noisy environment** is actually a **ally**.
- The Metrics of Success:
Increased **coherence time**,
understand **transport** of quantum
excitations/information.



The Toy (Minimal, solvable) Model

$H = H_S + H_{SE}$
 (The von Neumann Style)



$d\mathbf{E} = -\mathbf{A} \cdot \mathbf{E}dt + \mathbf{B} \cdot d\mathbf{W}$ (SDE: Ornstein-Uhlenbeck)
 $\Sigma dt\delta(t - t') = d\mathbf{W} \otimes d\mathbf{W}$ (correlation matrix)

The Toy (Minimal, solvable) Model

$$H = H_S + H_{SE}$$

(The von Neumann Style)

Unitary Evolution

Liouville-von Neumann Equation

$$\frac{d\rho_{SE}}{dt} = -i[H, \rho_{SE}] + \sum_i D_i(\rho_{SE})$$

Dissipation superoperator

$$D_i(\rho_{SE}) = L_i \rho_{SE} L_i^\dagger - \frac{1}{2} \{L_i^\dagger L_i, \rho_{SE}\} + L'_i \rho_{SE} L'^{\dagger}_i - \frac{1}{2} \{L'^{\dagger}_i L'_i, \rho_{SE}\}$$

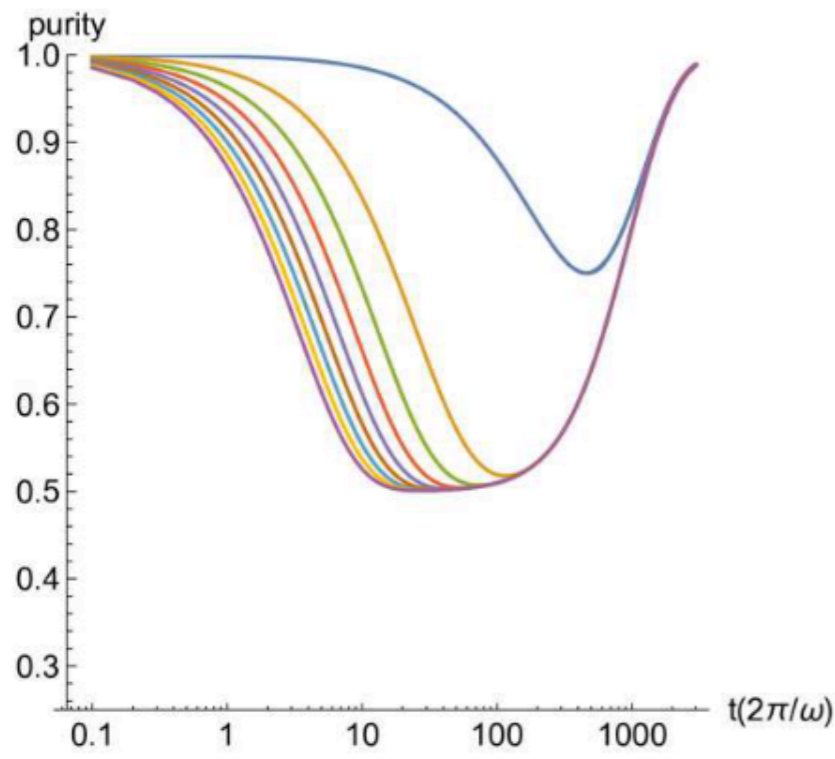
- $\xi = 0 \implies$ Uncorrelated Noise
- $\xi = -1 \implies$ Anti-correlated Noise
- $\xi = +1 \implies$ Correlated Noise

$$L_{\pm} = \sqrt{\frac{1 \pm \xi}{2}} (L_1 \pm L_2)$$

$$L'_{\pm} = \sqrt{\frac{1 \pm \xi}{2}} (L'_1 \pm L'_2)$$

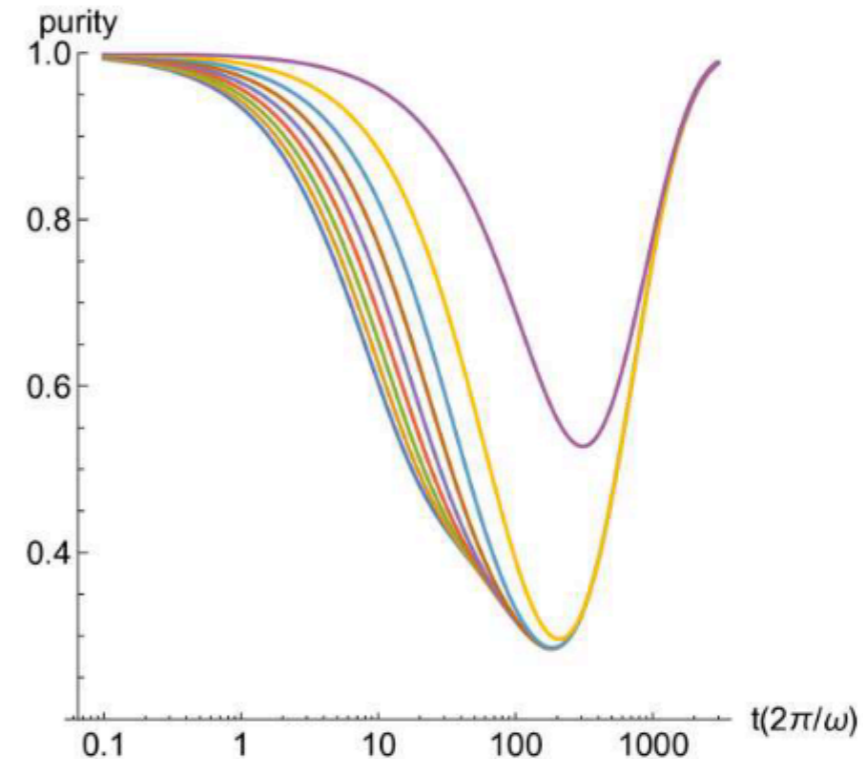
The correlation parameter!
 $-1 \leq \xi \leq +1$

We see it, but wait...



- $\xi = -1.$
- $\xi = -0.75$
- $\xi = -0.5$
- $\xi = -0.25$
- $\xi = 0$
- $\xi = 0.25$
- $\xi = 0.5$
- $\xi = 0.75$
- $\xi = 1.$

$$Tr(\rho^2)$$



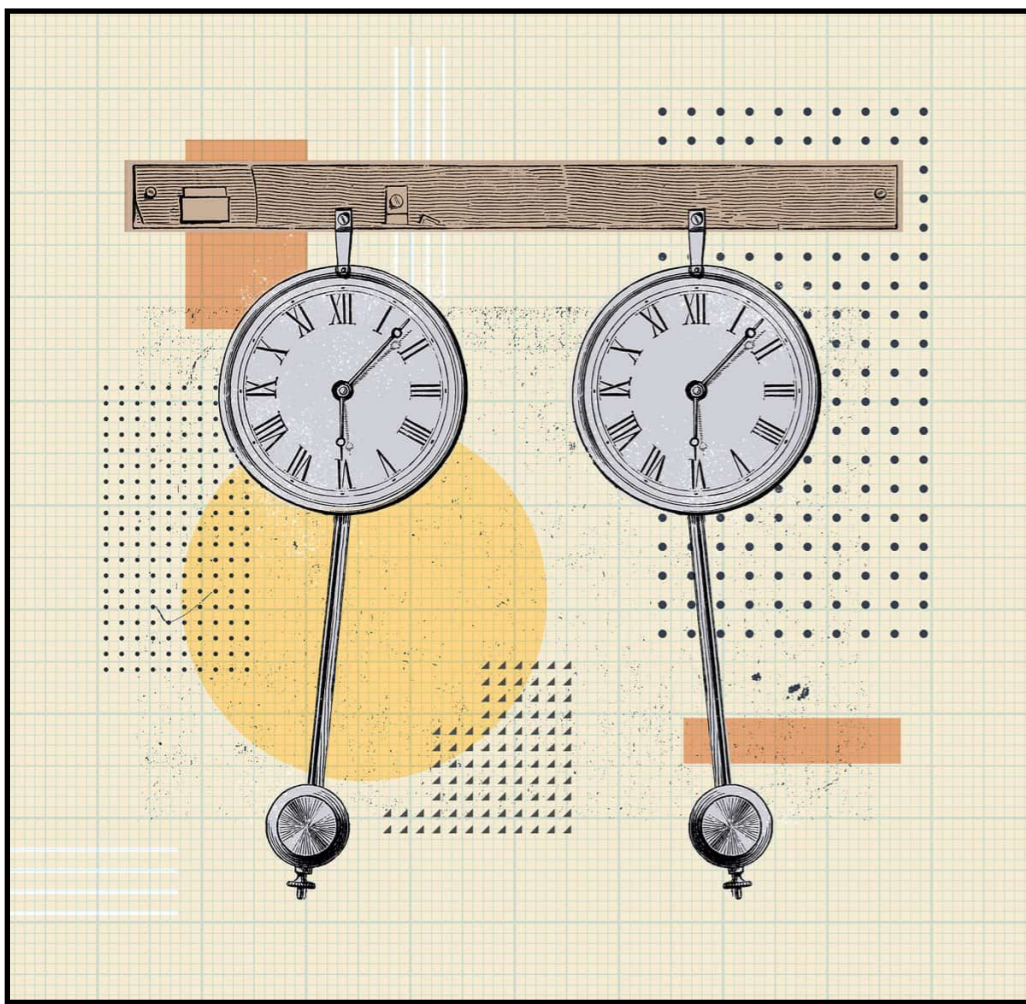
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- $\xi = -0.25$
- $\xi = 0$
- $\xi = 0.25$
- $\xi = 0.5$
- $\xi = 0.75$
- $\xi = 1.$

$$\Phi^+ = \frac{1}{2}(00 + 11)$$

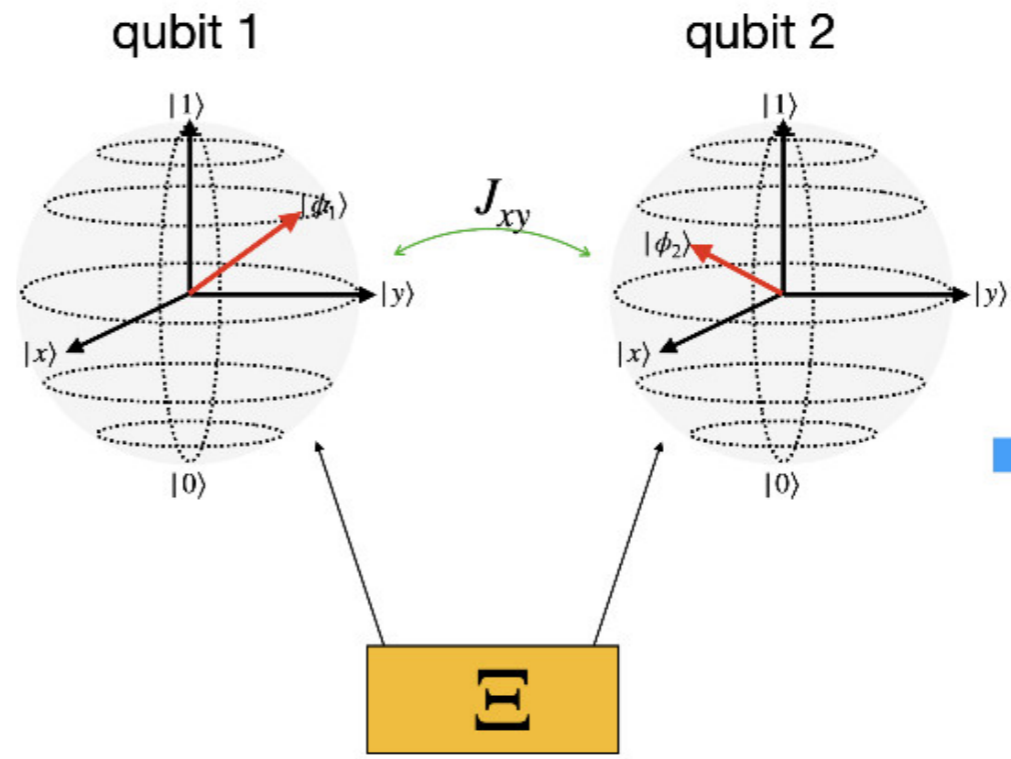
Increased coherence time

$$\Psi^+ = \frac{1}{2}(01 + 10)$$

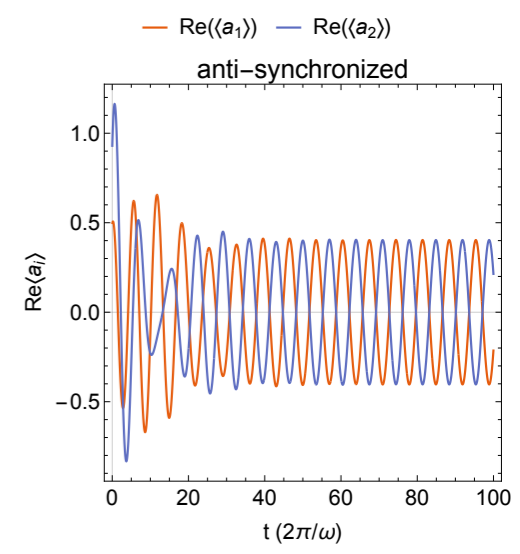
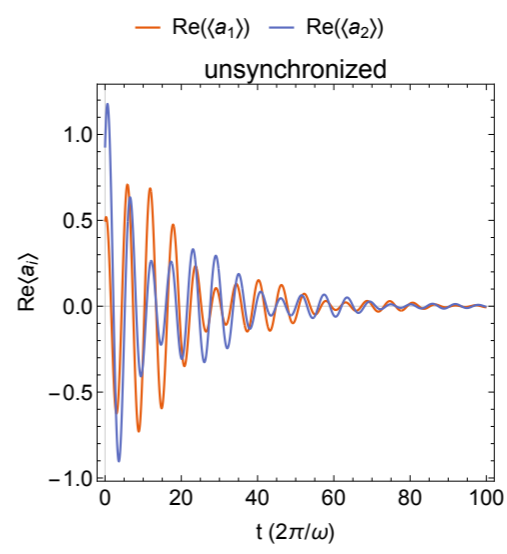
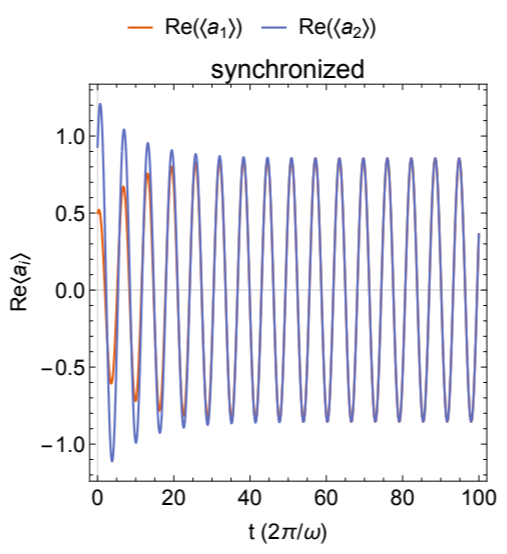
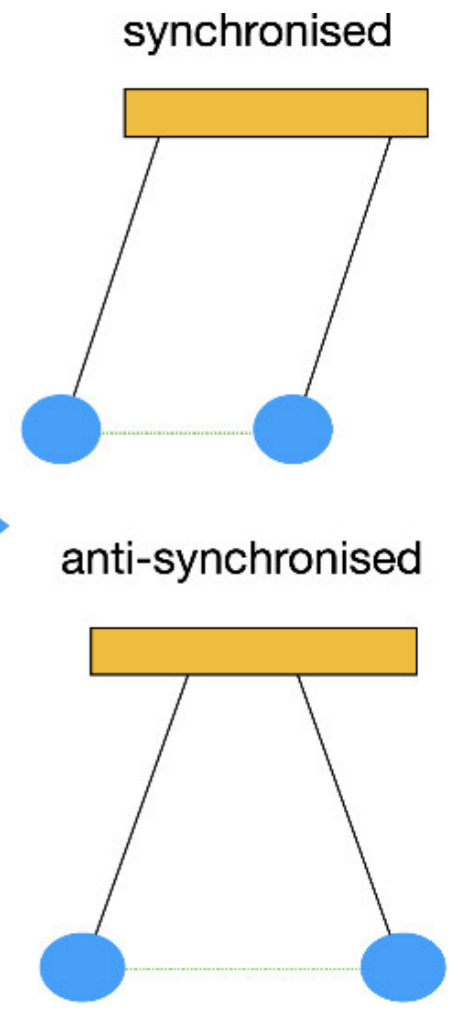
Lets Re-search



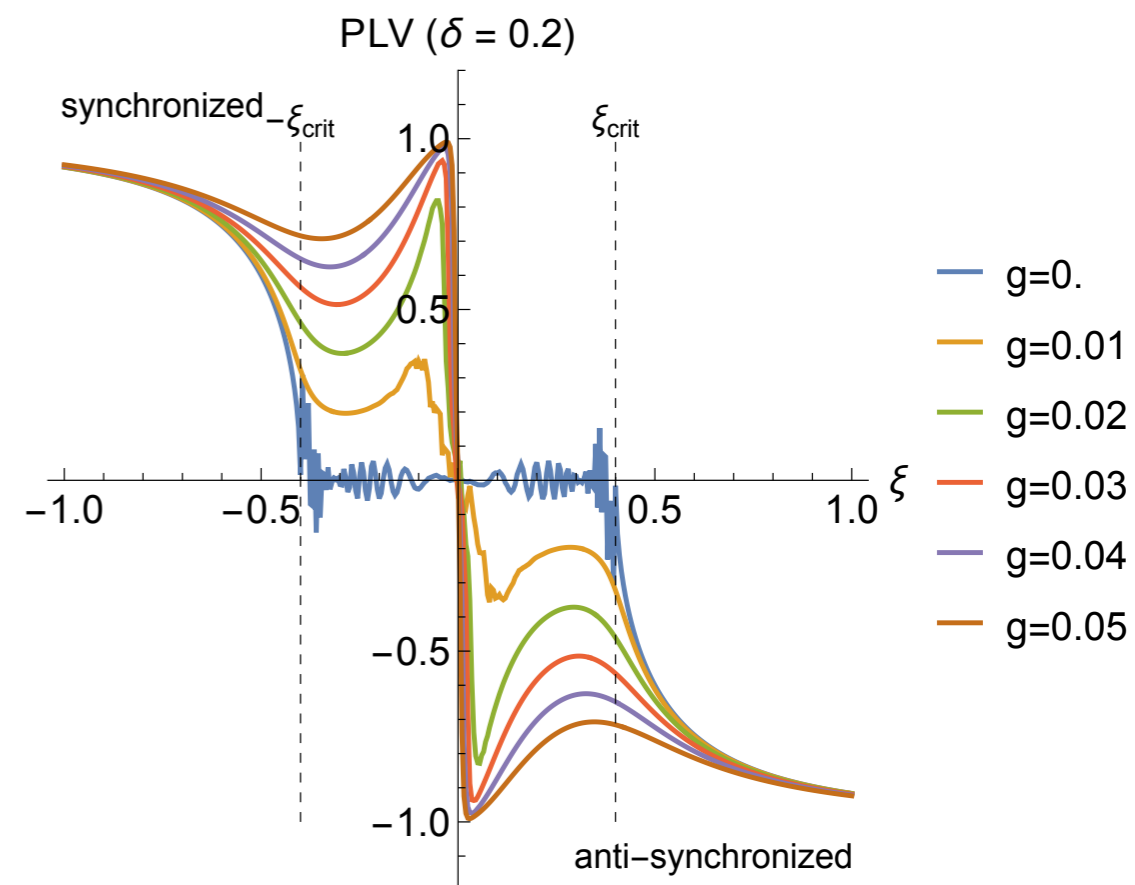
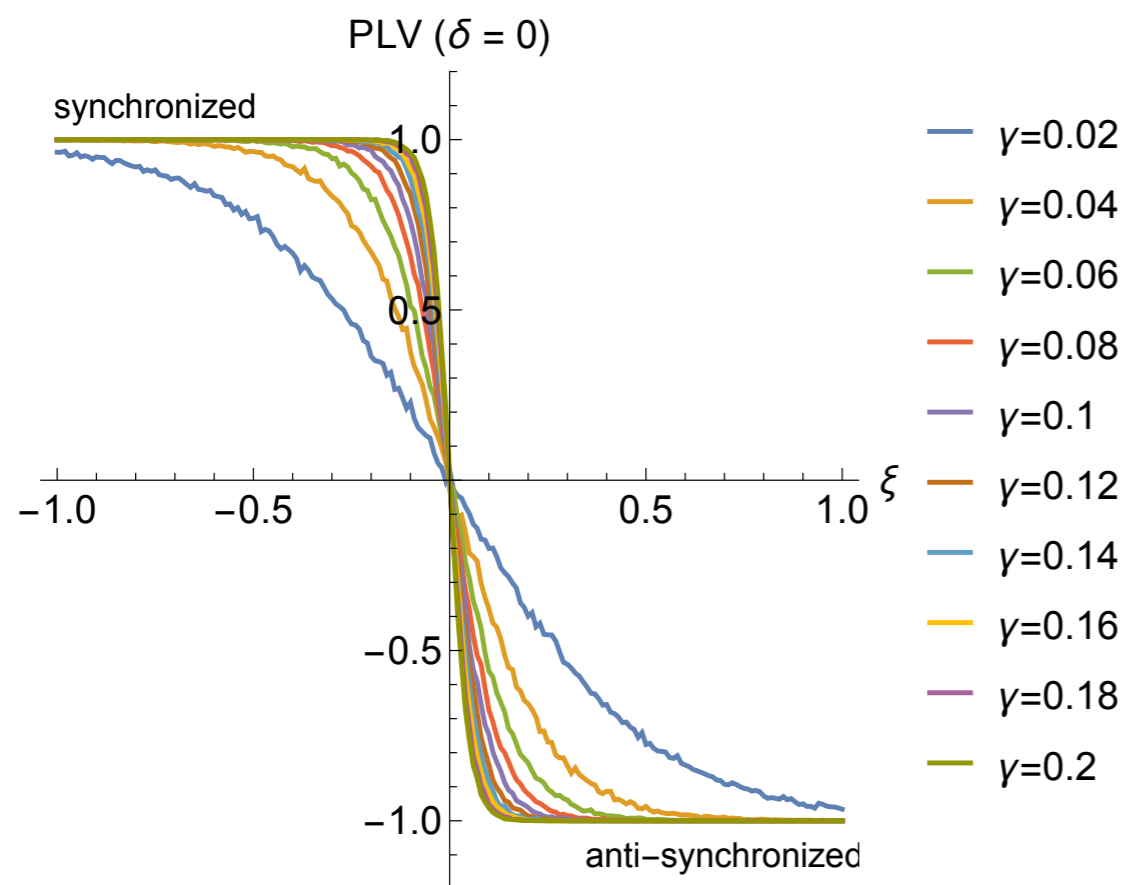
<https://physicsworld.com/a/the-secret-of-the-synchronized-pendulums/>



correlated common environment



Oh really?



Oh that's cool!

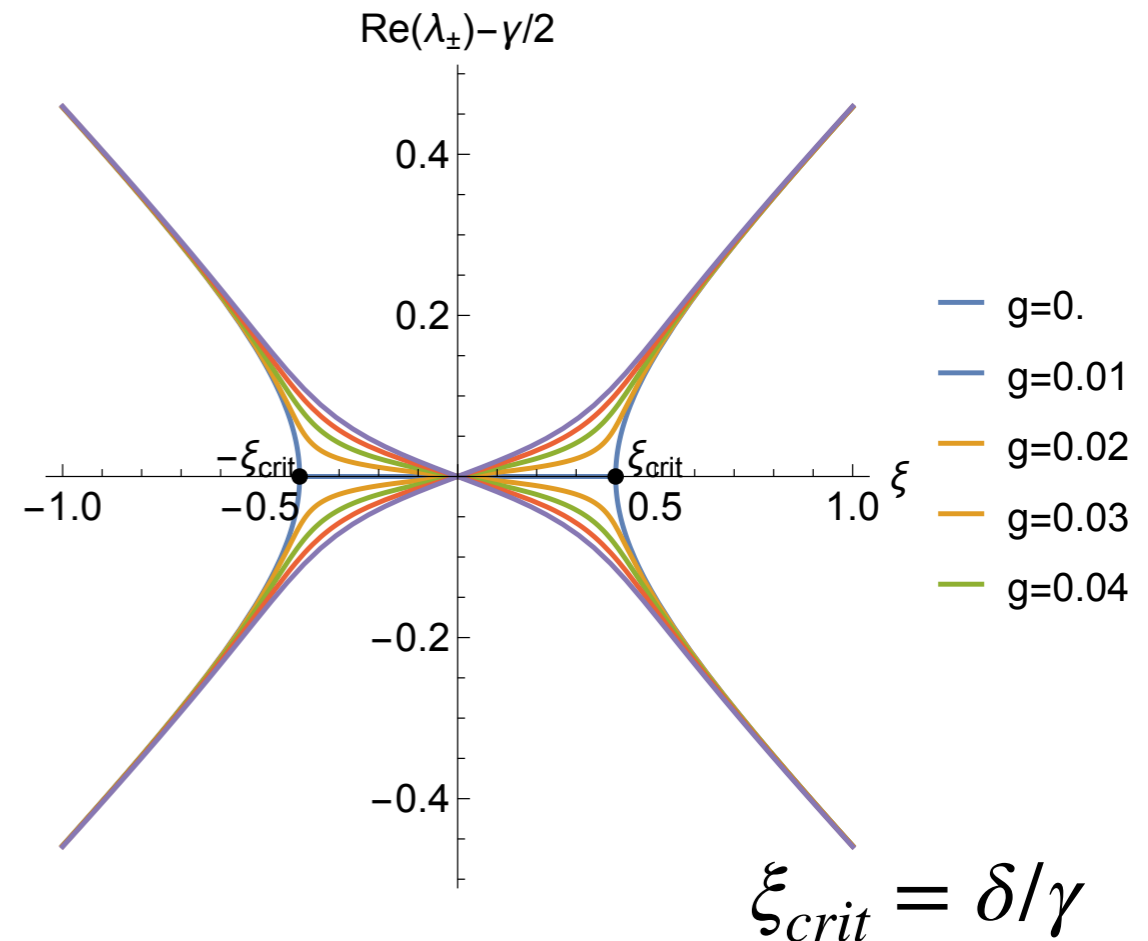
A novel and universal result.

\mathcal{PT} symmetry breaking.

Decoherence free subspaces

What does the future hold?

- System on different topologies.
- Larger system size.
- ξ is emergent in finite dimensional systems.

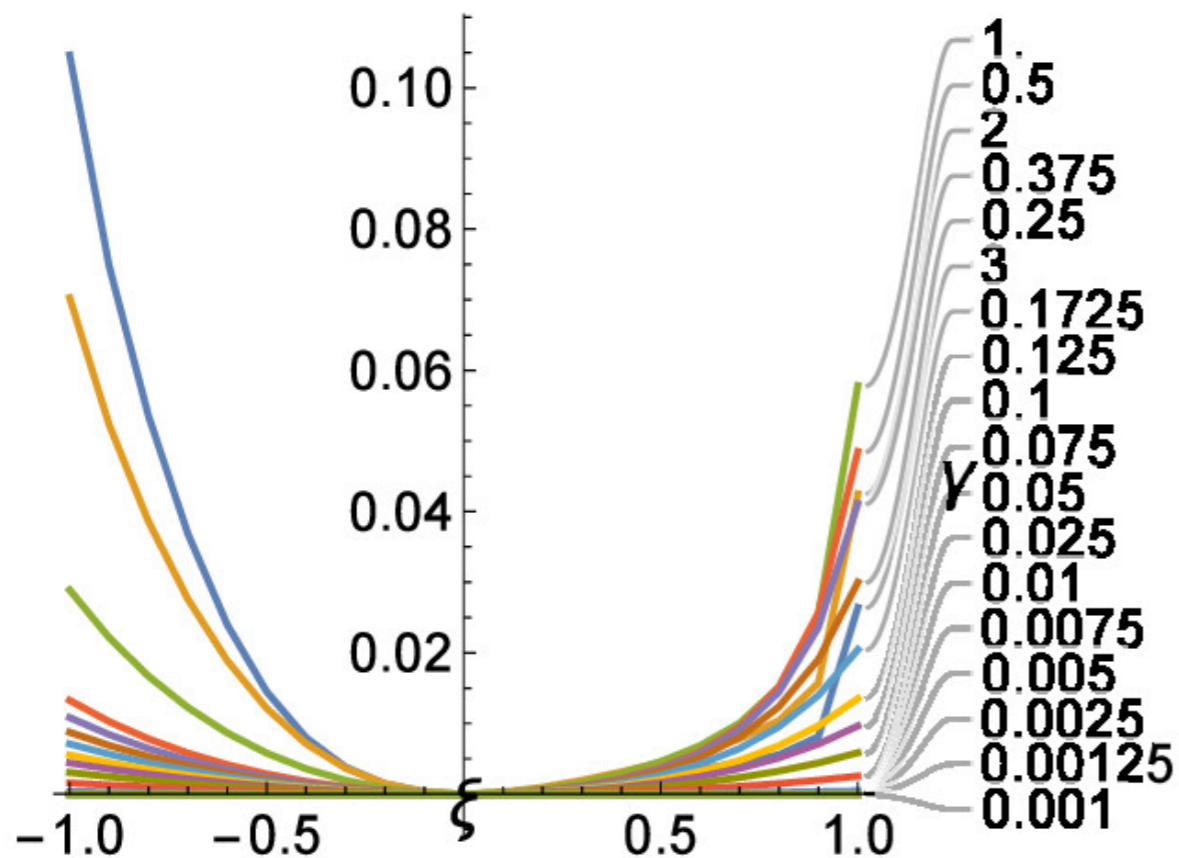


We did it, again! 😊

$$\rho_{SE} \xrightarrow{\text{Long-time}} \text{diag}(\rho_{SE}) = \begin{pmatrix} ii & \times & \times \dots & \times \\ \times & ii & \times \dots & \times \\ \times & \times & ii \dots & \times \\ \times & \times & \times \dots & ii \end{pmatrix}$$

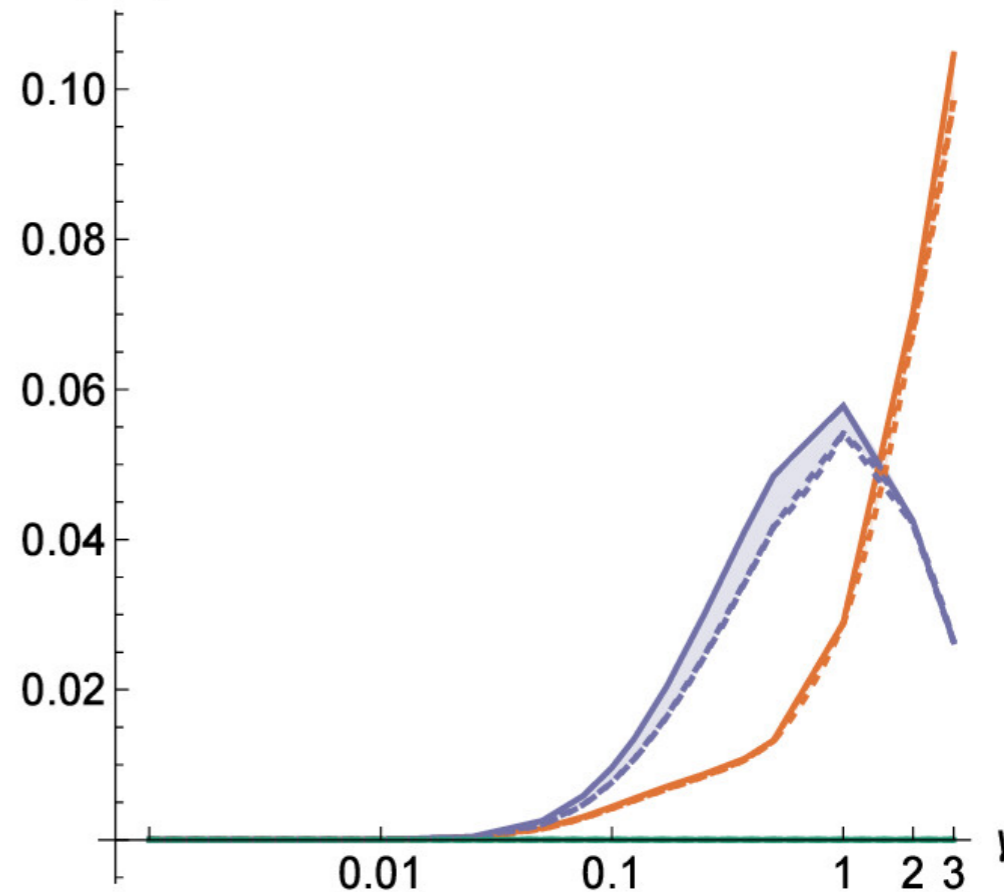
$$J_{xy}=0$$

$$I(A:B) = S_A + S_B - S_{AB}$$



$$J_{xy}=0$$

$$I(A|B) = \text{Classical Correlations} + \text{Quantum Correlations}$$



Decoherence (kind of) helped us.

Oh that's way cooler!

A novel and universal result.

$$D_i(\rho_{SE}) = L_i \rho_{SE} L_i^\dagger - \frac{1}{2} \{L_i^\dagger L_i, \rho_{SE}\} + L'_i \rho_{SE} L'^{\dagger}_i - \frac{1}{2} \{L'^{\dagger}_i L'_i, \rho_{SE}\}$$

$$L_{\pm} = \sqrt{\frac{1 \pm \xi}{2}} L_1 \pm L_2 \quad L'_{\pm} = \sqrt{\frac{1 \pm \xi}{2}} L'_1 \pm L'_2$$

- $\xi = 0 \implies$ Cross dissipation channels vanish.
- $\xi = -1 \implies$ Symmetric terms vanish.
- $\xi = +1 \implies$ Anti-symmetric terms vanish.

Quantum Noise-Cancelling Headphones

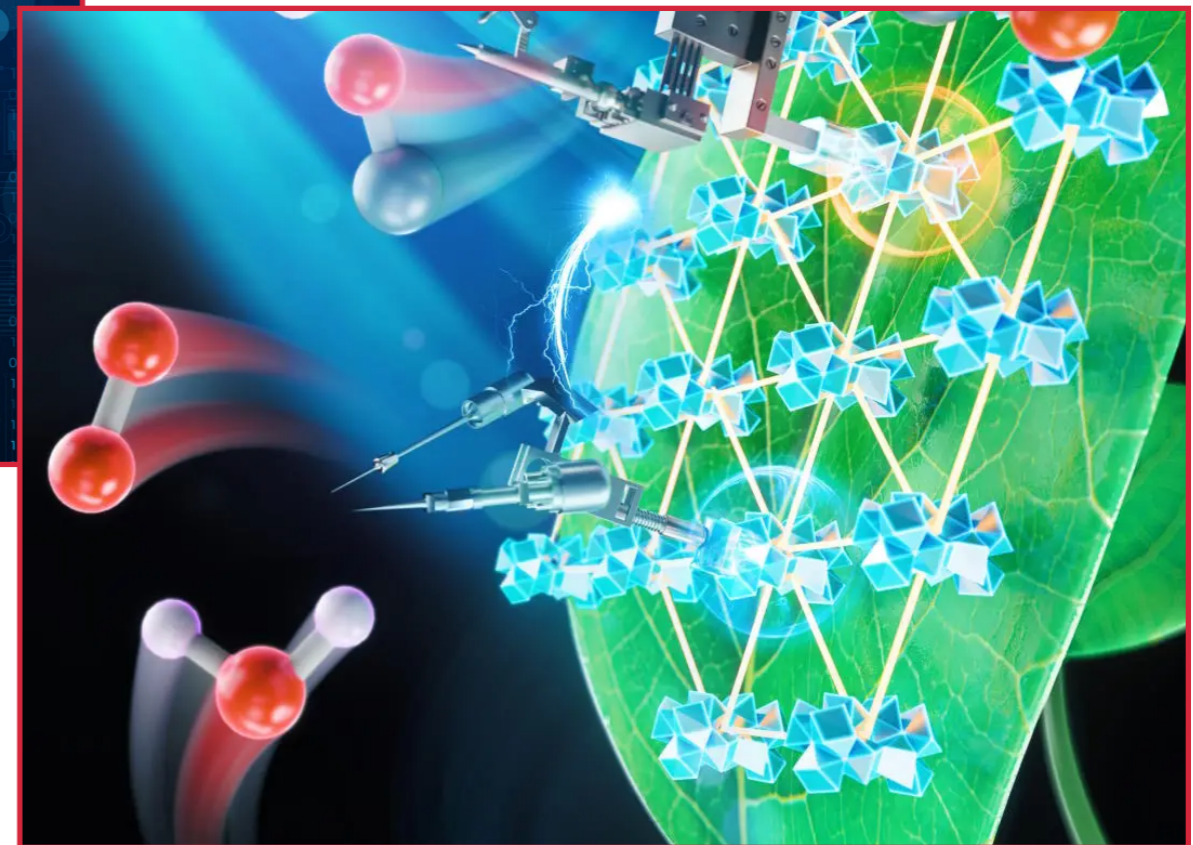


Real-world impact

<https://www.einfochips.com/blog/quantum-computing-in-artificial-intelligence-around-the-corner/>



“It is still an unending source of surprise for me how a few scribbles on a blackboard or on a piece of paper can change the course of human affairs.”
- Stanislaw Ulam



Thank you for your attention!

Peter Allen/University of Chicago

Back Up

$$\text{PLV} = \lim_{t_2 \rightarrow \infty} \frac{1}{t_2 - t_1} \text{Re} \left\langle \int_{t_1}^{t_2} e^{i(\phi_1(t) - \phi_2(t))} dt \right\rangle$$

$$D_{\Pi}(A | B) = I(A : B) - J_{\Pi}(A | B)$$

$$J_{\Pi}(A | B) = S(\rho_A) - S_{\Pi}(A | B)$$

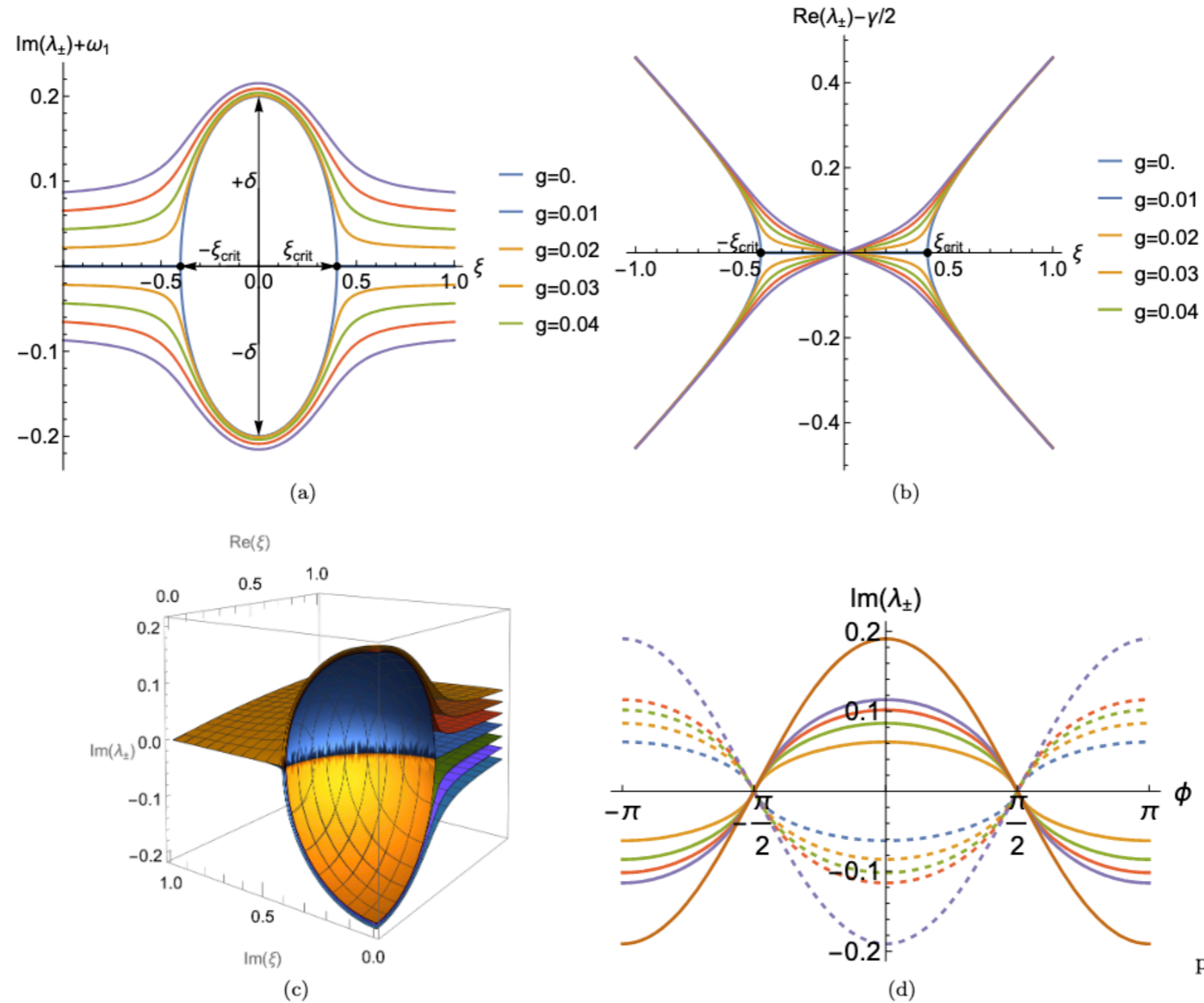


FIG. 1. Imaginary (a) and Real (b) components of the eigenvalues of the non-Hermitian dynamical matrix in Eq. 19 for various parametric values: $\delta = \omega_1 - \omega_2 = 0.2$, $\gamma = 0.5$ and increasing values of exchange coupling $g = 0$ to $g = 0.04$. The system enters the spontaneous synchronized domain above the critical coupling $\xi_{crit} = \delta/\gamma$. (c) Analytical continuation of (b) onto the complex ξ -plane for showing the presence of a branch cut starting at $\pm i\xi_{crit}$. (d) Slice through (c) along $|\xi_{crit}|$ passing through the exceptional points at $\phi = \pm\pi/2$ for $g > 0$. Solid: $\text{Im}(\lambda_+)$. Dashed: $\text{Im}(\lambda_-)$.